

## In Zeeman effect

**Paschen Back effect** :- Due to spin and orbital motion of  $e^-$  the external M.F. is weak as compared to internal field. When strength of M.F. is increased the sep. b/w Zeeman components increase until they become greater than separation b/w multiplet fine structure component. Therefore anomalous Zeeman effect changes over to like a normal Zeeman pattern. This phenomena is called Paschen Back effect.

When external field  $B$  becomes stronger as compared with internal field, the magnetic coupling b/w  $\vec{J}$  and  $B$  exceeds that between  $\vec{L}$  and  $\vec{S}$  and precession of  $\vec{J}$  about  $B$  becomes faster than  $\vec{L}$  &  $\vec{S}$  about  $\vec{J}$ . Under these condition the coupling b/w  $\vec{L}$  and  $\vec{S}$  is partially broken down and  $\vec{J}$  is no longer mixed in magnitude. As field  $B$  is further increased  $\vec{L}$  and  $\vec{S}$  start precessing independently about  $B$  with quantised component ( $L_z$  and  $S_z$ ) along field direction. The magnitude of these components are  $M_L \frac{h}{2\pi}$  and  $M_S \frac{h}{2\pi}$  resp.

$$M_L = L, L-1, \dots, -L$$

$$M_S = S, S-1, \dots, -S$$

By Larmour theorem

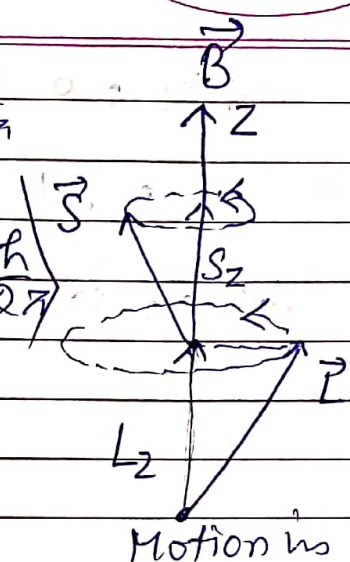
$$W_L = \frac{eB}{2mc}, \quad W_S = \frac{2e}{2mc} B$$

$$\Delta E_L = \omega_L L_z = \frac{eB}{2mc} \cdot \frac{M_L h}{2\pi}$$

$$\Delta E_S = \omega_S S_z = \frac{2eB}{2mc} \cdot \frac{M_S h}{2\pi}$$

$$\Delta E_2 = \Delta E_L + \Delta E_S$$

$$= (M_L + 2M_S) \frac{eh}{4\pi mc} B$$



shifts in wave number unit P.B. Effect

$$\Delta T = \frac{\Delta E}{hc} = (M_L + 2M_S) \frac{eB}{4\pi mc^2}$$

$$= (M_L + 2M_S) L'$$

Each field free level is split into  $(2L+1)(2S+1)$  magnetic levels.

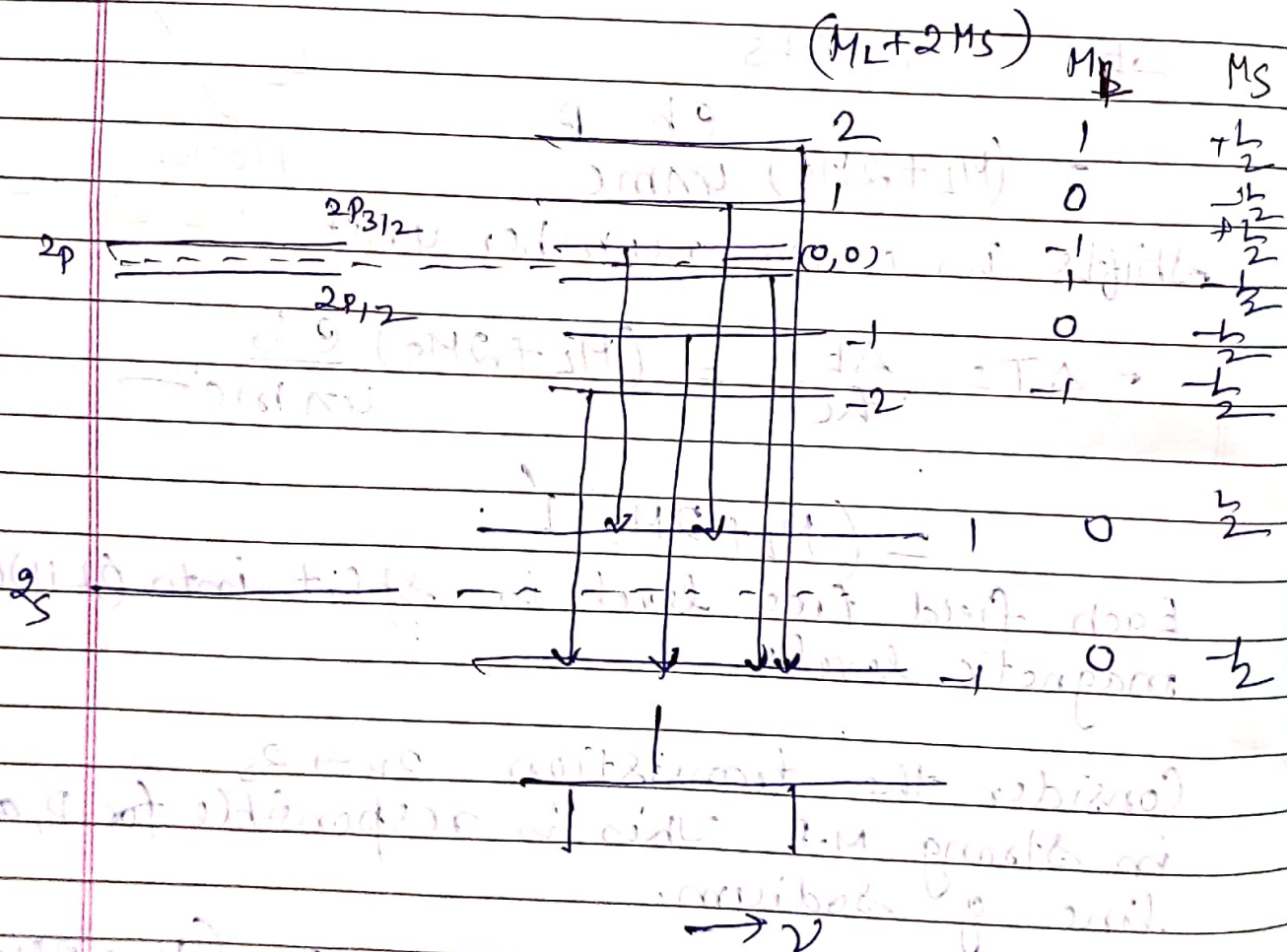
Consider the transition  $2p \rightarrow 2s$  in strong M.F. This is responsible for  $D_1$  and  $D_2$  line of sodium.

Term	No. of field levels	$M_L$	$M_S$	$(M_L + 2M_S)$
$2p$ $L=1, S=\frac{1}{2}$	6	1	$\frac{1}{2}, -\frac{1}{2}$	2, 0
		0	$\frac{1}{2}, -\frac{1}{2}$	1, -1
		-1	" "	0, -2
$2s$ $L=0, S=\frac{1}{2}$	2	0	" "	1, -1



$\Delta M_L = 0$  Component polarised || to field  
 $= \pm 1$   $L^2$

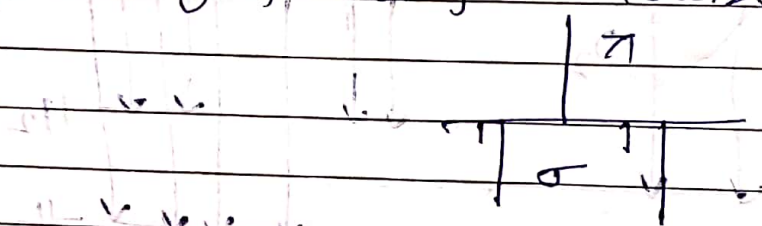
$\Delta M_S = 0$



**Spin orbit Correction:** In practice the residual spin orbit coupling changes the relative energies of the components of different terms. This can be allowed for by adding a small coupling term  $a M_L M_S$  in the expression for magnetic interaction energy which now becomes

$$-AT = (M_L + 2M_S) L' + a M_L M_S$$

Each of two  $\sigma$  components of normal triplet splits into narrow doublet, triplet accord to original field free transition was a doublet, triplet. In the present field free transition  $2p \rightarrow 2s$ , each  $\sigma$  component splits into a doublet with separation just  $2/3$  that of field free doublet.



The P.B. effect has experimentally been obs. for very low multiplets only such as Li doublet having field free separation of  $34 \text{ cm}^{-1}$ . The reason is that the effect occurs when magnetic splitting exceeds the fine structure splitting for field  $43000 \text{ Gauss}$ .